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Title:

Coupling fluid flow, heat transfer and thermal denaturation-aggregation of beta-lactoglobulin using an Eulerian/Lagrangian approach

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1 1. Introduction

2

3 Thermal processing techniques have been applied to improve quality and 4 safety of food products and to extend their shelf life (Wang and Sun, 2006). In the dairy 5 industry, billions of liters of milk are processed every year for pasteurization or sterilization. 6 Fouling of heat exchangers is an issue because it reduces heat transfer efficiency and increases pressure drop, hence affecting the economy of processing plants. As a result of 7 8 fouling, there is the possibility of deterioration in the product quality because the process 9 fluid cannot be heated up to the required temperature. Although representing only about 5 % 10 of the milk solids, whey proteins account for more than 50 % of the fouling deposits occurring at temperatures below 110 °C (Bansal and Chen, 2006). The understanding of 11 fouling mechanisms has progressed since the 1970's, which include not only the aggregation 12 13 of beta-lactoglobulin with other protein molecules but also the influence of aggregates on heat transfer between walls and the process fluid. 14

15 The structure of whey proteins can be significantly altered depending on the duration 16 and severity of heat treatment. With increasing temperatures whey proteins unfold and can 17 aggregate, either themselves or with other species. Aggregates including whey proteins are relevant to human nutrition (de Wit, 1998) and health (Navarra et al., 2007). Among other 18 19 applications in the food industry, whey protein aggregates with particle sizes ranging between 20 0.1 and 20 micrometers have been incorporated as successful fat replacers into low-fat stirred 21 yoghurts, because they provide sensory sensations close to those associated with the presence 22 of fat milk (Janhoj et al., 2006). Modeling of whey protein aggregation induced by heat 23 treatment can therefore provide assistance in understanding biophysical processes as well as 24 in designing processing units.

25 This study is devoted to the numerical modeling of the evolution of a suspension of 26 whey proteins under continuous thermal treatment. The beta-lactoglobulin is taken as an 27 example among the whey proteins, and for the sake of simplicity it is taken alone in the 28 suspension. It is the most abundant protein found in the whey fraction of cow milk, and also the dominant whey protein in heat-induced fouling (Bansal and Chen, 2006). Its structure and 29 30 properties have been studied since the 1930's with the help of different biochemical 31 techniques (Sawyer and Kontopidis, 2000). Denaturation-aggregation of beta-lactoglobulin 32 occurs following different stages: i) the dissociation of native dimers existing under 33 physiological conditions into native monomers, when the temperature increases above 40 °C; ii) the reversible partial unfolding of monomers and the formation of a thermally induced 34 molten globule state when the temperature increases above 60 °C; and iii) depending on 35 operating conditions, the occurrence of irreversible intra-molecular interactions resulting in 36 37 the formation of aggregates (Tolkach and Kulozik, 2007). An overview about the thermal behavior of the beta-lactoglobulin at temperatures up to 150 °C, including the influence of 38 pH, has been presented by de Wit (2009). 39

40 Reaction kinetic approaches have been applied for studying the thermal denaturation-41 aggregation of the beta-lactoglobulin. Such a strategy can provide information related to the main stages of the proteins evolution under heat treatment, namely the concentration of 42 43 native monomers and that of partially unfolded monomers, and then the aggregates 44 concentration (Tolkach and Kulozik, 2007). Nevertheless, kinetic descriptions cannot provide 45 information regarding the evolution of particle sizes along the process unit. According to 46 Bansal and Chen (2006), efforts remain to be accomplished in assessing the role of different 47 types of aggregates on the fouling deposit formation, as well as in generalizing mitigation 48 strategies. The evolution of particle sizes along the process unit can also be relevant to the sensory and rheological properties of products like voghurt (Torres et al., 2011). 49

50 The evolution of a suspension of beta-lactoglobulin under continuous thermal 51 treatment involves the occurrence of fluid flow, heat transfer and product transformation. 52 These processes are coupled: on one hand, the product transformation depends on the 53 operation conditions (velocity, shear and temperature fields); on the other hand, the liquid 54 food thermal and rheological properties are themselves modified as a result of the product 55 transformation. Such a two-way dependence is put in evidence in Figure 1. The suspension flows along a tubular heat exchanger, which is represented by an axially-symmetric 56 57 rectangular domain; the inlet and outlet are represented by the left and right sides of the 58 domain, and the suspension is heated from the wall (bottom of domain).

In Figure 1a, it is assumed that fluid parcels move according to a plug-flow velocity 59 field, and that they are exposed to a temperature field depending only on the distance from 60 the exchanger inlet. Such a one-dimensional (1D) model, neglecting the radial dependence of 61 62 both velocity and temperature, represents the simplest way to characterize the path to be followed by fluid parcels along a heat exchanger. At the inlet the suspension is assumed to 63 contain monomers only. Upon exposure to heat the temperature of suspension can reach the 64 65 threshold value for unfolding; after that, unfolded monomers can aggregate. At the outlet, the suspension contains remaining monomers and particles of different sizes resulting from 66 aggregation. In Figure 1a, the resulting particle size distribution is characterized by a 67 qualitative width A. 68

However, as shown in Figure 1b, one cannot neglect the fact that fluid parcels move slower near heating walls. Such fluid parcels undergo dynamical and thermal (and hence kinetic) histories in the process unit that are different from those associated with parcels running far from walls. As a consequence, the extension of unfolding and aggregation depends no longer on the distance from the exchanger inlet only, but on those histories experienced by each fluid parcel. The extension of the product transformation is therefore

3/29

75 expected to increase towards the wall. Different size distributions can be obtained at the 76 outlet, depending on the radial position (see Figure 1b); as a result, the bulk size distribution can be broader than that obtained with the 1D model (compare widths A and B). The 77 78 formation of larger particles allows an increase of the volume associated with the suspension 79 solid fraction, which in turn can modify the suspension viscosity. Figure 1b puts in evidence 80 the role played by dynamical and thermal histories experienced by fluid parcels, showing that both the velocity and the temperature depend on the radial distance. We argue that naive 81 82 assumptions about fluid flow and heat transfer (like the 1D model, Figure 1a) can provide a 83 biased picture regarding particle sizes at the heat exchanger outlet.

In order to realistically represent the denaturation-aggregation of beta-lactoglobulin 84 85 under thermal treatment, we propose a numerical approach that combines computational fluid dynamics (CFD) for solving the fluid-flow and heat-transfer coupled problem, with a 86 87 population balance equation (PBE) for evaluating particle aggregation. CFD has become an important tool for research and education in a number of subjects including food science and 88 89 technology (Norton and Sun, 2006). For instance, CFD has been employed in looking for the 90 optimal design of plate heat exchangers conceived for milk processing (Grijspeerdt et al., 91 2003). The PBE is a common model used to represent the evolution of particulates and 92 dispersed-phase systems when nucleation, growth, aggregation or fragmentation (breakage) 93 of particles occur (Marchisio and Fox, 2005). For instance, the PBE has been adapted in 94 studying the breakage of whey protein precipitates under selected flow conditions (Zumaeta et al., 2007; Bas et al., 2009). Simulation strategies combining CFD and PBE approaches 95 96 have became feasible with the increase of computation facilities, allowing a more detailed 97 analysis of the particle, droplet or bubble's distribution along the fluid evolution (Marchisio 98 and Fox, 2005; Bhole et al., 2008). In the scope of food science and technology, strategies

99 combining CFD and PBE have been implemented in studying the ice crystallization of
100 sucrose solutions in scraped surface freezers (Lian *et al.*, 2006).

101 In this exploratory study, our aim is to demonstrate that a more realistic representation 102 of the thermal denaturation-aggregation of whey proteins, and particularly of the aggregation 103 of particles under consideration, can be reached with an approach able to take into account 104 the different dynamical and thermal histories associated with fluid parcels in the processing 105 unit. The present approach combines the respective advantages of Eulerian and Lagrangian 106 descriptions of fluid flow, focusing on the coupled representation of fluid flow, heat transfer 107 and whey protein denaturation-aggregation processes. Fluid flow and heat transfer are solved through the finite element method in the Eulerian frame, while the product transformation is 108 evaluated along Lagrangian trajectories associated with representative fluid parcels. Whey 109 protein unfolding is represented by a reaction kinetics of order 1.5, while particle aggregation 110 111 is assessed after discretization of the PBE. The impact of particle sizes on the suspension viscosity is estimated, allowing the update of the previous solution for fluid flow and heat 112 transfer. The coupled problem is solved by iterating successive Lagrangian and Eulerian steps 113 up to a selected level of convergence. 114

115 The article is organized as follows. Section 2 presents how the processes of fluid flow, 116 heat transfer, proteins unfolding and particle aggregation are modeled. Section 3 describes 117 how the coupled problem is solved; particular attention is devoted to the solution of the PBE 118 after its discretization, and to the coupling between the PBE solution and CFD results. 119 Section 4 discusses relevant results from our Eulerian-Lagrangian approach, including a 120 comparison with those obtained by neglecting the radial dependence of velocity and 121 temperature fields along the heat exchanger (1D model). Finally, Section 5 provides a 122 summary and indicates future work.

123

5/29

124 **2. Modeling**

125

126	In this study, the suspension of beta-lactoglobulin is represented as a single-phase
127	system. The coupled problem of fluid flow and heat transfer processes is solved on the whole
128	domain using an Eulerian representation while the protein unfolding and particle aggregation
129	are evaluated on a set of fluid parcels whose Lagrangian trajectories depends on the fluid
130	velocity. Such a strategy allows the evaluation of transformation processes that cannot be
131	easily coupled with CFD. This is the particular case of the evaluation of particle aggregation
132	through the PBE approach by employing class or pivot methods, which represent the particle
133	size distribution with the help of a number of size classes.
134	
135	2.1. Transformation model
136	
137	2.1.1 Case study
138	
139	This work deals with the aggregation of beta-lactoglobulin under typical heat treatment. This
140	is achieved by assuming a protein concentration of 3.2 g.L ⁻¹ , which is a typical value found in
141	cow milk (e.g., Swaisgood, 1995), in water solution at pH 6.8 with 0.1 M NaCl added.
142	Choosing such values ensure that beta-lactoglobulin is under about the same electrochemical
143	environment as in milk. Data for beta-lacoglobulin unfolding and aggregation were obtained
144	for these conditions from experimental work (Aymard et al., 1996; Elofsson et al., 1996; Verheul
145	et al., 1998).
146	In order to test the modeling approach in industrial-like conditions, the coupled fluid flow -
147	heat transfer - transformation model was implemented inside a tubular heat exchanger with
148	length 1 m and radius 5 mm, under a flow rate (10 L.h ⁻¹) and a heating rate (10 kW.m ⁻²)

149	which lead to a temperature increase from 50 $^{\circ}\mathrm{C}$ (pre-heating temperature) to a maximum
150	value of about 110 °C. Such a treatment allows a significant level of unfolding-aggregation
151	and could be achieved on a pilot plant for further comparison with experimental data.
152	
153	2.1.2. Protein unfolding
154	
155	The denaturation-aggregation of beta-lactoglobulin has been modeled using a reaction
156	kinetics (Anema and McKenna, 1996). The native protein concentration C_{nat} (mol.m ⁻³) at
157	time <i>t</i> is given by:
158	

159
$$\frac{DC_{nat}\{t\}}{Dt} = -k_n C_{nat}^n\{t\}$$
(1)

160

161 where *n* is the reaction order and k_n is the rate constant of the reaction $((\text{mol.m}^{-3})^{(1-n)}.\text{s}^{-1})$. The 162 reaction order *n* may express the complexity of the unfolding sub-reactions, rather than the 163 formal order of a classical chemical reaction; in the case of beta-lactoglobulin, it usually 164 takes values between 1 and 2, commonly 1.5 (Tolkach and Kulozik, 2007). The relationship 165 between the rate constant and the temperature of the reaction is given by the Arrhenius 166 equation:

167

168
$$k_n = k_0 \exp\left(\frac{-E_a}{RT}\right)$$
(2)

169

where k_0 is the frequency factor, E_a the activation energy (J.mol⁻¹), *R* the universal gas constant (J.mol⁻¹.K⁻¹) and *T* the temperature (K).

Experimental work on the thermal denaturation-aggregation of the beta-lactoglobulin has put in evidence a marked change in temperature dependence of the reaction rate constant at a threshold of about 85-90 °C (e.g., Lyster, 1970). As a matter of fact, such a behavior results from the occurrence of two regimes: below the temperature threshold, the reaction is limited by the unfolding sub-reactions; above that value, the reaction is limited by the aggregation process (Tolkach and Kulozik, 2007).

In this study the thermal denaturation-aggregation of beta-lactoglobulin is represented as a two-step process. Firstly, a reaction kinetics is assumed in estimating the concentration of native proteins which undergo partial reversible unfolding; the temperature is assumed to be the only limiting factor. Secondly, the population balance equation is solved in order to estimate the final concentration of the particles which reach, via aggregation, a number of prescribed sizes. Limiting factors of this second step are the concentration of unfolded proteins, the temperature and the shear rate.

A reaction kinetics of order 1.5 is hereafter assumed in representing the thermal unfolding of the beta-lactoglobulin. Table 1 summarizes values of $k_{1.5}$ obtained by Anema and McKenna (1996) for two variants of beta-lactoglobulin under selected temperatures in the range 70-85 °C. Looking for a more representative set of results, an unique Arrhenius equation is assumed to characterize both variants. We estimated the frequency factor and the activation energy by fitting the data presented on Table 1 as a whole:

191

192
$$k_0 = 8.27 \ 10^{42} \ (\text{mol.m}^{-3})^{(-0.5)} \text{.s}^{-1}$$
 (3)

193

194 $E_a = 3.04 \ 10^5 \ \text{J.mol}^{-1}$

195

(4)

196 These values are employed into equations (1) and (2) independently on the value assumed by 197 the temperature as computed by the numerical approach. For temperatures above 85 °C the 198 extrapolation of the Arrhenius equation can lead to under- or overestimate the unfolding rate, 199 but unfolding is then no more the limiting factor. Indeed, in our case study, unfolding times 200 rapidly decrease with the temperature (~1 min at 85 °C, ~1 s at 100 °C); unfolded protein 201 molecules are immediately engaged in the irreversible aggregation process which becomes 202 the limiting mechanism. JUS

203

2.1.3. Particle aggregation 204

205

206 Once unfolded, the whey protein monomers are assumed to aggregate according to the population balance equation proposed by Smoluchowski (1917). This kind of model has been 207 208 used for a wide range of applications, including the aggregation of whey proteins (Wu et al., 209 2005; Péron et al.; 2007). After assuming no breakage of aggregates, the evolution of the number of particles containing k monomers (N_k) can be expressed as: 210

211

212
$$\frac{DN_{I}\{t\}}{Dt} = k_{I.5} C_{nat}^{1.5}\{t\} N_{A} - N_{I}\{t\} \sum_{\ell=I}^{M} Q_{I,\ell} N_{\ell}\{t\}$$
(5a)

213
$$\frac{DN_k\{t\}}{Dt} = \frac{1}{2} \sum_{\ell=1}^{k-1} Q_{k-\ell,\ell} N_\ell\{t\} N_{k-\ell}\{t\} - N_k\{t\} \sum_{\ell=1}^M Q_{k,\ell} N_\ell\{t\}$$
(5b)

214

215 in the cases k = 1 and k > 1, respectively, where $N_1(t)$ is the number of unfolded monomers 216 at time t, M the maximum number of monomers in an aggregate and N_A the Avogadro 217 number. The kernel of aggregation $Q_{k,l}$ depends on the local temperature, on the continuous

218 phase viscosity (η_w) , on the shear rate $(\dot{\gamma})$ (Pa.s), and finally on the radii a_k and a_l of the 219 aggregates under consideration:

220

221
$$Q_{k\ell} = \frac{2}{3} \frac{k_B}{W} \frac{T}{\eta_w} \frac{(a_k + a_\ell)^2}{a_k a_\ell} + \frac{4}{3} \frac{\dot{\gamma}}{W} (a_k + a_\ell)^3$$

222

where k_B is the Boltzmann constant and W the Fuchs stability factor. The radius a_k of an aggregate containing k monomers can be estimated from the monomer radius a_k and the aggregates fractal dimension (*D*):

226

227
$$a_k = a_0 k^{1/D}$$

(7)

228

The fractal dimension describes the microstructure of the aggregates (Bremer et al., 1989), allowing to link the number of monomers in an aggregate with its size. This parameter depends on the monomer / continuous phase couple as well as on the operating conditions (Byrne et al., 2002).

233 In this exploratory study we focus the attention on the aggregation of beta-234 lactoglobulin molecules. Under the considered conditions (pH=6.8 and 0.1 M NaCl), beta-235 lactoglobulin monomers exhibit a globular hydrodynamic radius (a_0) of about 2 nm (Aymard 236 et al., 1996). Under quite similar conditions, comparisons between measurements and 237 numerical PBE solutions have suggested the value W=91 for the Fuchs stability factor 238 (Elofsson et al. 1996). Still under those conditions, the combination of permeability, 239 rheology, and aggregation kinetics measurements on whey protein gels has indicate the value 240 D = 2.4 for the aggregates fractal dimension (Verheul *et al.*, 1998).

241	Taking into account the concentration of 3.2 g.L ⁻¹ , the initial suspension contains
242	$1.0531 \ 10^{23}$ monomers per cubic meter. For a maximum total volume fraction of 0.6,
243	corresponding to the random packing, the largest aggregates should have a radius of about 10
244	μ m and should contain about $M = 10^9$ monomers.
245	
246	2.2. Fluid flow and heat transfer models
247	
248	2.2.1. Mass, momentum and energy conservation
249	
250	The conservation equations for mass, momentum, and energy can be written as:
251	
252	$\nabla \cdot (\rho \mathbf{v}) = 0 \tag{8}$
253	
254	$\rho\left(\mathbf{v}\cdot\nabla\right)\mathbf{v} = \nabla\left(-p\cdot\mathbf{I} + \eta\left(\nabla\mathbf{v} + (\nabla\mathbf{v})^T\right)\right) $ (9)
255	
256	$\rho C_P \left(\mathbf{v} \cdot \nabla \right) T = \nabla \cdot \left(\lambda \nabla T \right) \tag{10}$
257	
258	Coupled phenomena are illustrated inside a tubular heat exchanger with length 1 m
259	and radius 5 mm, which is represented through a two-dimension axial-symmetric domain. A

fully developed parabolic flow profile was assumed at the inlet. The flow rate was fixed at 10 liters per hour (mean velocity 0.0351 m.s^{-1}). A uniform flux of 10 kW.m⁻² is applied on the heating wall, and the temperature associated with the suspension is assumed to be 50 °C at the inlet. These conditions enable temperatures of about 110 °C at the outlet near the heating wall. Such a temperature enables a relevant degree of protein unfolding and particle

aggregation (Tolkach and Kulozik, 2007; de Wit, 2009). According to the Reynolds number,

the flow regime is laminar. Boundary conditions are summarized in Figure 2.

267

268 2.2.2. Suspension viscosity

269

The liquid food product under consideration is an aqueous suspension of whey proteins. The apparent viscosity η associated with the suspension is obtained by multiplying the continuous phase (pure water) viscosity η_W and the relative viscosity η_R which is a function of the transformation state. Pure water viscosity depends on the local temperature, and such an influence can be evaluated through a reference formulation (IAPWS, 2008):

275

276
$$\eta_W = 10^{-6} \left(280.68 \left(\frac{T}{300} \right)^{-1.9} + 511.45 \left(\frac{T}{300} \right)^{-7.7} + 61.131 \left(\frac{T}{300} \right)^{-19.6} + 0.45903 \left(\frac{T}{300} \right)^{-40} \right)$$
277 (11)

278

279 Various approximations have been proposed for estimating the relative viscosity as a 280 function of the volume fraction (Φ) occupied by the particles under consideration. We have 281 retained the model proposed by Thomas (1965),

282

283
$$\frac{\eta}{\eta_W} = \eta_R = 1 + 2.5 \ \Phi + 10.05 \ \Phi^2 + 0.00273 \ exp(16.6 \ \Phi)$$
 (12)

284

because it provides an increase of relative viscosity with the particle volume fraction which is intermediate to those obtained with other approximations. The volume fraction is calculated from the third moment of the aggregates size distribution:

288

(13)

289
$$\Phi = \frac{4\pi}{3} \sum_{k=1}^{M} N_k \{t\} a_k^3$$

290

3. Problem solving

292

3.1. Population balance equation resolution

294

Once unfolded, the beta-lactoglobulin is subjected to aggregation, represented by the 295 population balance equation. Its resolution by equation (5) involves the calculation of the 296 population of a very large number of aggregates size, as an aggregate can contains up to 10^9 297 298 monomers. In order to reduce the computational cost, only a limited number of pivots, 299 indexed by i or j, are here considered. This approach has been used by Batterham et al. (1981) and Hounslow et al. (1988); the latter computed the size distribution of kidney stones 300 by considering only the particles whose volume is on a regular grid described by $V_{i+1} / V_i = 2$. 301 This approach can be applied to proteins aggregation by considering the monomers number 302 of a particle instead of its volume. Here, a monomer number grid is used with $n_{i+1} / n_i = 2$, 303 where n_i is the number of monomers in the aggregates of pivot *i*. The aggregates of the pivots 304 *i* hence contain $n_i = 2^{i-1}$ monomers. Figure 3 displays an example of population balance and 305 306 its repartition on such a pivot grid.

307 As expected, the conservation of quantities such as the total mass or the total number 308 of particles may be problematic. To solve this issue, (Kumar and Ramkrishna, 1996) showed 309 that some conditions can lead to the conservation of several quantities using only a reduced 310 set of pivots.

For instance, in the case of the aggregation of particles of size k = 2 and l = 8 (as shown on Figure 4), the resulting particle contains 10 monomers and thus, is not among the calculated pivots. In order to take its presence into account, the number of particles in the two

314 nearest pivots ($k = 8 \Leftrightarrow i = 4$ and $k = 16 \Leftrightarrow j = 5$) are increased by A for pivot i and B for 315 pivot *j*. In order to conserve the total number of monomers and of particles in this particular 316 case, A and B must satisfy: for the conservation of particles number: 317 A + B = 1318 for the conservation of monomers number: $A \cdot 8 + B \cdot 16 = 10$ • This gives A = 3/4 and B = 1/4 that are the variations on pivots 4 and 5. 319 Regarding the conservation of the total number of monomers and particles on a 320 geometric grid $(n_{i+1} = 2 n_i)$, the proposed solution is equivalent of the earlier proposition of 321 (Hounslow et al., 1988) that describes the time evolution of the number of particle at the 322 NA pivot *i*, denoted N^i , by: 323 271

325 for *i* = 1:

326
$$\frac{dN_{I}\{t\}}{dt} = k_{I.5} C_{nat}^{I.5}\{t\} N_{A} - N_{I}\{t\} \sum_{j=I}^{M} Q_{I,j} N_{j}\{t\}$$
(14a)

327 for $l < i \leq M$:

$$\frac{dN_{i}\{t\}}{dt} = N_{i-1}\{t\} \sum_{j=1}^{i-2} 2^{j-i+1} Q_{i-1,j} N_{j}\{t\} + \frac{1}{2} Q_{i-1,j-1} N_{i-1}^{2}\{t\} - N_{i}\{t\} \sum_{j=1}^{i-1} 2^{j-i} Q_{i,j} N_{j}\{t\} - N_{i}\{t\} \sum_{j=i}^{M} Q_{i,j} N_{j}\{t\}$$
(14b)

329

330 For the general case of equation 14b, the first term accounts for the creation of aggregates on the pivot *i* by aggregation of aggregates from pivot i-1 with smaller ones, whereas the second 331 332 term accounts for the creation of aggregates on pivots *i* by aggregation of two aggregates from pivot i-1. The last two terms described the destruction of aggregates on pivot i by 333 334 aggregation with either smaller (third term) or larger (last term) aggregates. The number of 335 pivots is limited to M = 30, corresponding to aggregates of radius about 10 µm.

336

337 3.2. Coupling of fluid flow, heat transfer and product transformation by an Eulerian338 Lagrangian approach

339

A schematic diagram of the approach is presented in Figure 5. The liquid food is assumed flowing in a tubular heat exchanger, which is represented by a rectangular domain. The dashed red and the bold black lines indicate the axis of symmetry and the heating wall, while the left and right sides represent the domain inlet and outlet respectively. Fluid parcels running nearer to the heating wall move slower than those at the axis; therefore dynamical and thermal histories differ accordingly, leading to different transformation state evolution along their trajectories.

The present study combines the respective advantages of Eulerian and Lagrangian descriptions of fluid flow. Fluid flow and heat transfer are solved through the finite element method in the Eulerian frame, while the product transformation is evaluated along Lagrangian trajectories associated with representative fluid parcels.

Before the first iteration, a preliminary modeling of fluid flow and heat transfer is needed. It provides a first estimate of the velocity, shear rate and temperature fields along the Eulerian frame. Because this modeling step is performed before any product transformation, the viscosity is here assumed to be the one of the continuous phase (pure water). Eulerian modeling of fluid flow and heat transfer is solved by using a mesh constituted of 100 rectangles along the domain's radius and 2000 rectangles along the length.

The first step of a given iteration is the estimation of Lagrangian trajectories associated with a number of representative fluid parcels. Each trajectory is defined by a number of positions, and each of them is associated with a given time since the fluid parcel release at the domain inlet. The estimation of trajectory positions from the Eulerian velocity

361 field is performed with the help of a routine implementing a fourth-order Runge-Kutta 362 method (Abramowitz and Stegun, 1972). After a number of sensitivity tests, we found 363 satisfactory the application of 200 trajectories with about 200 positions on each.

In the second step, Lagrangian trajectories are employed in the reconstruction of dynamical and thermal histories along the domain. Each history is described by a series of positions, times, and the values associated with the shear rate and temperature. These values are estimated from the last available Eulerian respective fields.

Next, the shear rate and temperature histories are employed, along every trajectory, in evaluating the protein unfolding degree (eqs. 1 and 2) and in solving the PBE (eqs. 6 and 14). This is the key step regarding the flexibility of the Eulerian-Lagrangian approach proposed: any available methodology can be employed in evaluating the degree of transformation of the liquid food product along the process unit. The availability of the particle size distribution enables the evaluation of the volume fraction associated with the solid fraction (eq. 13), which is then used for estimating the suspension viscosity (eq. 12).

375 Updated suspension viscosity values have to be informed to the Eulerian model, as a 376 previous task towards an updated solution for the fluid flow and heat transfer problem. 377 Difficulties at this step come from the fact that Lagrangian trajectories provide updated 378 viscosity values at a number of positions which are irregularly distributed along the Eulerian 379 frame. The Delaunay triangulation was retained as strategy. The application of constraints 380 relying successive positions along a same trajectory proves useful. As a result, the suspension 381 viscosity values available at Lagrangian positions are successfully interpolated into a 382 suspension viscosity field needed at the nodes constituting the Eulerian frame.

383 Before the following iteration, a new Eulerian modeling of fluid flow and heat 384 transfer is performed. It provides updated fields for the velocity, shear rate and temperature. 385 Because this modeling step is performed by taking into account an improved suspension

386 viscosity field, it is expected that these velocity, shear rate and temperature fields represent a 387 somewhat better picture of the fluid flow and heat transfer conditions under the occurrence of 388 product transformation.

The coupled problem is solved by iterating the successive Lagrangian and Eulerian steps up to a selected level of convergence. The solution is assumed to be reached when successive viscosity fields exhibit a maximum difference smaller than 0.25 %.

This algorithm was implemented through a main program and a set of functions developed in the programming language MATLAB. Eulerian modeling of fluid flow and heat transfer was solved with the help of the COMSOL Multiphysics 3.5a simulation package, using the UMFPACK solver. The implementation of both tasks, including the adoption of other mathematical solvers, seems feasible in other programming languages.

397

398 3.3. One-dimensional model

399

The transformation of liquid food product under heat treatment has been often represented in the literature by considering average temperature evolution along the exchanger and by assuming plug flow. In this case, both the velocity and temperature fields are assumed to be independent on the distance from the heating wall (e.g., Grijspeerdt et al. (2004)). Such an one-dimensional (1D) model is included in the present study for the sake of comparison with a more elaborate approach, in which no assumptions are performed regarding the velocity and temperature fields.

407 Coherently with the application of the Eulerian-Lagrangian approach, the evaluation 408 of the thermal denaturation-aggregation of beta-lactoglobulin requires the unfolding rate 409 constant, (eq. 2) and the aggregation kernel (eq. 6), whose computation asks for local 410 estimates of temperature and shear rate. In the scope of the 1D model, bulk values are used.

411 Assuming that the liquid food product is exposed to an uniform heat flux density Q $(W.m^{-2})$, the heat balance for a tubular exchanger can be expressed as: 412 413 $\rho \bar{v} \pi R^2 C_P d\bar{T} = 2\pi R Q dz$ 414 (15)415 where ρ (kg.m⁻³) is the product density, $\overline{\nu}$ (m.s⁻¹) the plug-flow velocity, R (m) the radius of 416 the exchanger, C_P the specific heat capacity $(J.kg^{-1}.K^{-1})$ and \overline{T} (K) the bulk temperature. 417 The integration of equation (15) provides the bulk temperature as a function of the 418 distance from the exchanger inlet: 419 nA 420 $\overline{T}\{z\} = T_{inlet} + \frac{Q}{2 \rho \,\overline{v} \,R \,C_P} \,z$ 421 (16)422 A first-order estimate of the shear rate bulk value required for evaluating particle 423 aggregation in the 1D model may be obtained from the parabolic flow profile. At a given 424 425 radial distance r from the axis of symmetry of the tubular exchanger, the shear rate associated with parabolic flow can be expressed as: 426

427

$$428 \qquad \dot{\gamma}\{r\} = \frac{4\,\bar{\nu}}{R^2}\,r\tag{17}$$

429

430 Its bulk value is therefore given by:

431

432
$$\bar{\gamma} = \frac{\int \dot{\gamma}\{r\} v\{r\} dS}{\int v\{r\} dS} = \frac{32 \, \bar{\nu}}{15 \, R}$$
 (18)

434 Taking into account the geometry and boundary conditions defined above (R = 5 mm and \overline{v} $= 0.0354 \text{ m.s}^{-1}$), we obtain $\overline{\dot{\gamma}} = 15.1 \text{ s}^{-1}$. 435

436 After estimating bulk values for both the temperature and the shear rate, the 437 evaluation of the thermal denaturation-aggregation of beta-lactoglobulin can hence be performed by employing bulk values for both the temperature and the shear rate. For 438 MANG 439 instance, equation (1) can be rewritten as:

(19)

440

441
$$\frac{DC_{nat}\{t\}}{Dt} = \overline{v}\{z\}\frac{dC_{nat}\{z\}}{dt} = -k_n\{z\}C_{nat}^n\{z\}$$

442

- 4. Results 443
- 444
- 445 4.1. Fluid flow and temperature

446

Figure 6 shows the velocity, temperature, suspension viscosity and shear rate fields 447 resulting from the simulations described above. 448

449 The velocity field (Figure 6a) is assumed as a fully developed flow (parabolic profile) 450 at the domain inlet. The velocity at the exchanger axis (top of domain) firstly slightly 451 decreases along the first 30 % of the tube length, and then increases up to the outlet where it 452 reaches higher values than at the inlet. Near the heating wall, fluid parcels move slower. 453 Moreover, the region associated with relatively low velocities is broader near the outlet than 454 at the inlet.

455 Figure 6b shows that temperatures increase along the domain as a consequence of the 456 developing thermal boundary layer from the heating wall towards the exchanger center.

Figure 6c shows that the suspension viscosity increases near the heating wall along the domain, reaching values many times higher near the outlet than that observed at the inlet; moreover, the region associated with increasing suspension viscosity becomes progressively broader along the domain. The latter result is consistent with the decrease of the axial velocity near the heating wall (see Figure 6a).

462 At the domain inlet, the shear rate is proportional to the distance from the exchanger axis as a consequence of the fully-developed incoming flow. However, the shear rate field 463 464 displayed on the Figure 6d denotes 3 regions: weak values near the exchanger center, b) low 465 values at the wall after an initial increase, and c) between these two patterns, progressively increasing values in a ridge-like region that appears after about 20 % of the exchanger length 466 from the inlet. The first feature is consistent with the symmetry of the velocity profile at the 467 exchanger axis, while the second feature is consistent with the weak velocities occurring in 468 469 the region near the wall which becomes broader along the domain. The ridge-like pattern with higher shear rates is a consequence of previous features. 470

471

472 **4.2. Product transformation along trajectories**

473

The evolution of the product properties will be analysed along Lagrangian trajectories associated with fluid parcels released at the domain inlet. Each fluid parcel can be considered as a suspension droplet containing an ensemble of whey protein particles at similar thermodynamic and kinetic conditions.

Figure 6a displays the five trajectories hereafter taken into account, each of them associated with a given radial release positions at the domain inlet (r_0). These distances were chosen in order to sample different regions of the domain: firstly the exchanger axis, where the highest velocity as well as the minimum transformation are expected ($r_0 = 0$); then the

482 average radial position ($r_0 = 2.5$ mm); and finally three release positions progressively closer 483 to the heating wall ($r_0 = 4, 4.5, \text{ and } 4.75$ mm).

Figure 7a presents the temperature evolution along these five trajectories, the residence time values from 13.8 up to 174 seconds are indicated on the right. Fluid parcels move slower as they run closer to the wall, being therefore more exposed to heating. Coherently, outlet temperatures increase with r_0 .

The evolution of shear rate along these trajectories is displayed in Figure 6b. The shear rate is null on the axis and presents maximum values for the fluid parcels which are moving in the vicinity of the heating wall. The behavior of the two trajectories closest to the wall ($r_0 = 4.5$ and 4.75 mm) indicates first an increase, and then a decrease towards the outlet. The maximum values are associated with the ridge-like feature appearing in Figure 6d. Fluid parcels released near the middle of the domain ($r_0 = 2.5$ and 4 mm) are associated with simpler shear rate evolutions as they do not enter in the ridge-like region.

Temperature drives the unfolding kinetics (eqs. 1-2), and drives together with the shear rate the aggregation by modifying the aggregation kernel (eq. 6). As shown in Figure 7, these two key variables exhibit different evolutions along the domain from a trajectory to another. One can therefore expect different unfolding/aggregation histories along the domain under consideration, implying different whey protein particle size distributions at the outlet.

500 Figure 8 presents the evolution of the unfolding ratio -- that is the percentage of 501 proteins that have been unfolded; the volume fraction associated with whey protein 502 aggregates (Φ) and the Sauter mean diameter ($d_{3,2}$).

The liquid food product transformation is negligible for fluid parcels moving on the axis ($r_0 = 0$ mm) as well as for those released at the average radial position ($r_0 = 2.5$ mm). This is a consequence of the negligible whey protein unfolding because of the relatively low temperatures experienced along these two trajectories (see Figure 7a).

507 On the other hand, fluid parcels moving in the vicinity of the heating wall ($r_0 = 4.75$ 508 mm) experience almost complete unfolding after about 80% of the exchanger length as 509 displayed on Figure 8a. Such a high level of product transformation results from the fast 510 temperature increase over 70 °C. For $r_0 = 4.75$ mm, shear rate is also relatively high along the 511 first 60 % of the trajectory as depicted on Figure 7b; these conditions provide favorable 512 conditions to aggregation. This results in a large increase of the Sauter diameter (Figure 8b) 513 and of the volume fraction of the aggregates. The solid volume fraction increases up to 40 % 514 and the Sauter diameter becomes more than 2000 times larger than its value at the inlet. 515 Intermediate behaviors are exhibited by trajectories associated with fluid parcels released MA 516 between $r_0 = 2.5$ and 4.75 mm.

517

4.3. Viscosity evolution 518

519

The suspension viscosity is driven by the continuous phase (pure water) viscosity as 520 521 well as by the volume fraction of particles in suspension. Figure 9 shows that the suspension 522 viscosity can appreciably vary along the domain; moreover it can be very different from a 523 given trajectory to another.

524 Fluid parcels moving at the exchanger axis ($r_0 = 0$) and those released at the average 525 radial position ($r_0 = 2.5$ mm) are associated with negligible product transformation (see 526 Figure 8). The suspension viscosity decreases along these trajectories, indicating the role of 527 increasing temperature on the continuous phase viscosity.

On the other hand, fluid parcels moving closer to the heating wall ($r_0 = 4.5$ and 4.75 528 529 mm) experience both a meaningful heating and a high level of product transformation. As a 530 result, the suspension viscosity along these trajectories exhibits a two steps behavior (see Figure 9). On the first half of the exchanger, the suspension viscosity decreases (because 531

532 continuous phase viscosity is lower for higher temperatures). Then, the product is 533 transformed and aggregation occurs; it involves the increase of the suspension viscosity 534 because that the influence of the growing aggregates volume fraction overtakes the 535 temperature effect on the continuous phase.

536

537 **4.4. Product properties at the outlet**

538

After one meter of exchanger, the output product is subjected to a broad diversity due to the different time-temperature-shear rate histories. Figure 10a displays the cumulative volume distribution of the monomers and aggregates at the domain outlet for five selected trajectories.

For the two inner fluid parcels ($r_0 = 0$ and 2.5 mm), the aggregation is negligible, the monomers account for more than 98 % of the total volume. The trajectory starting at $r_0 = 4$ mm still contains a large number of monomers, but they account only for 6% of the solid volume.

For the fluid parcels starting from $r_0 = 4.5$ and 4.75 mm all the proteins are aggregated (no monomer remains). One can sum up that the closer the trajectory is to the heating wall, the higher is the median diameter.

Figure 10b displays the particle size distribution of the output bulk product. The output product described using the 2-D Eulerian-Lagrangian model is a blend of the different fluid parcels which a sample is displayed on Figure 10a. Consequently, a large number of monomers averagely remains, they account for 14 % of the solid volume The bulk median radius value is about 1 μ m with broad particle size distribution: 20 % of the solid volume is constituted by particles smaller than 0.2 μ m and 30 % by particles larger than 5 μ m.

556

557 **4.5. Comparison with 1D model results**

558

559	We have argued in the Introduction and shown previously that thermal and dynamic
560	histories can affect the transformation state of the liquid food product.
561	Table 2 presents the product properties at the outlet for both the 2D and the 1D model.
562	The first line displays the ranges of values assumed by selected variables at the domain outlet
563	(<i>r</i> varying from 0 up to $R = 5$ mm), as provided by our 2D simulations. Such ranges are very
564	large, clearly putting in evidence the variety of dynamic and thermal histories experienced by
565	fluid parcels.
566	The following lines in Table 2 present the bulk properties of the output product, that is
567	the properties of the fluid obtained by mixing all the outcoming fluid parcels. These values
568	are those of industrial process interest.
569	Bulk values are far from those associated with fluid parcels released at the average
570	radial position, as shown in the third line of Table 2. This 'average' trajectory ($r_0 = 2.5$ mm) is
571	thus not representative for the process.
572	Last line in Table 2 displays corresponding results provided by the 1-D model
573	described above. This model does not represent the variety of conditions experienced by the
574	liquid food product along the exchanger. It provides essentially the same results as those
575	obtained for fluid parcels released at the average radial position ($r_0 = 2.5$ mm). The results are
576	clearly different from the bulk estimated with the more realistic 2D coupled model.
577	
578	5. Conclusion
579	
580	Thermal denaturation of beta-lactoglobulin in a tubular exchanger has been modeled
581	using a 2D Eulerian/Lagrangian numerical scheme. By splitting the fluid flow and heat

24/29

transfer model from the transformation determination, a realistic population balance modelcan be used to estimate the proteins aggregation.

The results show the diversity of the output product properties depending on the timetemperature-shear rate histories of the different flow parts: in the region near the wall, the product is heavily heated for a long time, resulting in largely modified product; in the axial region, where the temperature is relatively low and the residence time is shorter, the product is weakly modified.

The described numerical approach takes into account the diversity of dynamic and thermal histories experienced by fluid parcels along different trajectories in contrary to the approach where the product is considered as a whole. Indeed, comparison with a simple 1D model shows that bulk properties of the outlet product cannot be predicted in this case without a 2D fully coupled approach.

The use of the Eulerian/Lagrangian approach allows the implementation of complex transformation models, as often necessary when dealing with food processing modeling, without the need of oversimplification or model rewriting to comply with the CFD solver structure.

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604

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697 List of tables

- 698
- Table 1: Chemical kinetics parameters from Anema and McKenna (1996).
- 700
- Table 2: Selected variables at the domain outlet.

	T [°C]	$[(mol.m^{-3})^{-0.5}.s^{-1}]$	R
β–lactoglob	ulin A 70 75 80 85	$ \begin{array}{r} 4.40 \ 10 \\ 2.31 \ 10^{-3} \\ 1.08 \ 10^{-2} \\ 2.68 \ 10^{-2} \end{array} $	
β–lactoglob	ulin B 70 75 80 85	$ \begin{array}{r} 3.60 \ 10^{-4} \\ 2.62 \ 10^{-3} \\ 1.46 \ 10^{-2} \\ 3.96 \ 10^{-2} \end{array} $	
	0 0		

	Unfolding ratio	Sauter diameter	particle volume	number concentration
	[%]	[m]	fraction [%]	[part/m ³]
2-D (range)	0 to 100 %	$2.0\ 10^{-9}$ to $8.6\ 10^{-6}$	0.35 to 53.62	$2.41 \ 10^{14} \text{ to } 1.053 \ 10^{23}$
2-D (bulk)	14.1	1.5 10 ⁻⁸	2.35	9.05 10 ²²
$2-D (r_0 = 2.5 \text{ mm})$	0.1	2.0 10-9	0.35	$1.052 \ 10^{23}$
1-D model	0.7	2.1 10-9	0.36	$1.046\ 10^{23}$

List of figures 1

- 2
- 3 Figure 1: Evolution of a suspension of beta-lactoglobulin under continuous heat treatment. (a)

4 simple 1D approach; (b) coupled 2 D approach.

5

6 Figure 2: Numerical domain and boundary condition.

7

8 Figure 3: Pivots representation.

9

US CRIF Figure 4: Particles repartition along pivots after aggregation. 10

11

Figure 5: Iterative procedure. 12

13

Figure 6: Velocity (a), temperature (b), viscosity (c), and shear rate (d) fields on the domain. 14

15 Domain inlet and outlet are on the left and right respectively, and the heating wall is on the

16 bottom. Five trajectories associated with fluid parcels released at the domain inlet are 17 illustrated with the axial velocity field (a).

18

19 Figure 7: Evolution of temperature (a) and shear rate (b) along fluid parcels trajectories.

20

21 Figure 8: Evolution of unfolding ratio (a), Sauter diameter (b), and aggregate volume fraction 22 (c) along the fluid parcels trajectories.

23

24 Figure 9: Evolution of suspension viscosity along the fluid parcel trajectories.

25

- 26 Figure 10: Cumulative volume distribution of monomers and aggregates at the outlet for: a)
- 27 five selected trajectories, b) bulk product.

Acception





Constant heat flux density $Q = 10 \, kW.m^{-2}$

Figure3











~

Figure8



Figure9





- We study β -lactoglobulin aggregation during continuous thermal treatment. •
- Aggregation is estimated along Lagrangian trajectories using sizes class method. •
- Transformation is coupled with fluid flow and heat transfer classically estimated. •
- Thermal and shear rate histories along the domain drive the aggregation magnitude. •
- Obtained particle sizes distribution shows product diversity unlike 1D model. •

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Title: Coupling fluid flow, heat transfer and thermal denaturation-aggregation of beta-lactoglobulin using an Eulerian/Lagrangian approach

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Abstract: The transformation of liquid food product under heat treatment is often represented by considering average temperature evolution along the exchanger and by assuming plug-flow. Our aim is to demonstrate that thermal denaturation-aggregation of whey proteins can be more realistically represented by taking into account the different dynamical and thermal histories associated with fluid parcels which progress more or less quickly, far or close to the heating wall, inside the processing unit. A numerical approach is proposed for evaluating the thermal denaturation-aggregation of whey proteins, combining computational fluid dynamics and the population balance equation. The approach is illustrated by the evolution of a suspension of beta-lactoglobulin. Fluid flow and heat transfer are solved through the finite element method in the Eulerian frame, while product transformation is evaluated along representative Lagrangian trajectories. Outlet bulk results show that the product reaches a higher level of transformation than in assuming plug-flow and radially-independent temperature.